## Mechanics of Thin-walled Structures

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1. Basic concepts of mechanics of structures
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b. Strain
c. Moment of inertia
i. First (static) moment of area (static)
ii. Second moment of area (inertia)
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## Aircraft Structures for Engineering Students

Fourth Edition


Stress, Strain, Hooke`s law, moments of area (inertia)

## State of stress



## State of stress

3D body is in equilibrum under action of externally applied forces $P$.


P - forces
O - point in the body
$\delta \mathrm{P}$ - resultant force
nn - plane which divides the body into two parts
$\delta \mathrm{A}$ - small area


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$$
\text { Stress }=\lim _{\delta A \rightarrow 0} \frac{\delta P}{\delta A}
$$

The direction of $\delta \mathrm{P}$ is not normal to the area $\delta \mathrm{A}$, in which case it is usual to resolve $\delta \mathrm{P}$ into two components:
$\delta P_{n}$ - normal to the plane
$\delta P_{s}$ - acting in the plane itself
Plane conatining $\delta \mathrm{P}$ is perpendicular to $\delta \mathrm{A}$. The stresses associated with the abovementioned components are:

Normal stress or direct stress

$$
\sigma=\lim _{\delta A \rightarrow 0} \frac{\delta P_{n}}{\delta A}
$$

Shear stress

$$
\tau=\lim _{\delta A \rightarrow 0} \frac{\delta P_{s}}{\delta A}
$$

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## State of stress

The resultant force $\delta \mathrm{P}$ may be resolved into the following components:

- one normal component
- two in-plane components of shear stress



## State of stress

What is the stress?
$\rightarrow$ vector
$\rightarrow$ tensor


## State of stress

What is the stress?
$\rightarrow$ vector
$\rightarrow$ tensor

The following elements must be specified to determine the stress:

- magnitude
- direction
- plane on which stress acts

$$
\text { Stress }=\lim _{\delta A \rightarrow 0} \frac{\delta P}{\delta A}
$$



## State of stress

What is the unit of stress?

$$
\left[\frac{N}{m^{2}}\right]=[P a]
$$



## State of stress

Normal and shear stresses:


Generally, except in cases of uniform stress, the normal and shear stresses on opposite faces of an element are not equal but differ by small amounts.

## State of stress

After solving, the equations of equilibrum can be written:

$$
\left.\begin{array}{l}
\frac{\partial \sigma_{x}}{\partial x}+\frac{\partial \tau_{x y}}{\partial y}+\frac{\partial \tau_{x z}}{\partial z}+X=0 \\
\frac{\partial \sigma_{y}}{\partial y}+\frac{\partial \tau_{y x}}{\partial x}+\frac{\partial \tau_{y z}}{\partial z}+Y=0 \\
\frac{\partial \sigma_{z}}{\partial z}+\frac{\partial \tau_{z x}}{\partial x}+\frac{\partial \tau_{z y}}{\partial y}+Z=0
\end{array}\right\}
$$

X, Y, Z - body forces coming from gravitational forces and inertia effects

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\frac{\partial \sigma_{z}}{\partial z}+\frac{\partial \tau_{z x}}{\partial x}+\frac{\partial \tau_{z y}}{\partial y}+Z=0
\end{array}\right\}
$$

X, Y, Z - body forces coming from gravitational forces and inertia effects
2D case - plane stress:

$$
\left.\begin{array}{l}
\frac{\partial \sigma_{x}}{\partial x}+\frac{\partial \tau_{x y}}{\partial y}+X=0 \\
\frac{\partial \sigma_{y}}{\partial y}+\frac{\partial \tau_{y x}}{\partial x}+Y=0
\end{array}\right\}
$$

$$
\begin{gathered}
\sigma_{z}=0 \\
\boldsymbol{\tau}_{\boldsymbol{x z}}=0 \\
\boldsymbol{\tau}_{\boldsymbol{y z}}=0
\end{gathered}
$$

## State of stress

## Mohr's circle of stress


(a)


## State of strain



## State of strain

Longitudinal or direct strains:
Shear strains:

$$
\left.\varepsilon=\lim _{L \rightarrow 0} \frac{\Delta L}{L} \quad \varepsilon_{x}=\frac{\partial u}{\partial x}, \begin{array}{rl}
\gamma_{x z} & =\frac{\partial w}{\partial x}+\frac{\partial u}{\partial z} \\
\varepsilon_{y} & =\frac{\partial v}{\partial y} \\
\varepsilon_{z} & =\frac{\partial w}{\partial z}
\end{array}\right\} \quad \gamma_{x y}=\frac{\partial v}{\partial x}+\frac{\partial u}{\partial y},
$$

## Stress - strain relations 1-D case

$$
\varepsilon_{x}=\frac{\sigma_{x}}{E} \quad \varepsilon_{y}=-v \frac{\sigma_{x}}{E} \quad \varepsilon_{z}=-v \frac{\sigma_{x}}{E}
$$

## Stress - strain relations 3-D case

$$
\left.\begin{array}{rl}
\varepsilon_{x} & =\frac{1}{E}\left[\sigma_{x}-v\left(\sigma_{y}+\sigma_{z}\right)\right] \\
\varepsilon_{y} & =\frac{1}{E}\left[\sigma_{y}-v\left(\sigma_{x}+\sigma_{z}\right)\right] \\
\varepsilon_{z} & =\frac{1}{E}\left[\sigma_{z}-v\left(\sigma_{x}+\sigma_{y}\right)\right]
\end{array}\right\}
$$

$$
\begin{aligned}
\sigma_{x} & =\frac{\nu E}{(1+v)(1-2 \nu)} e+\frac{E}{(1+\nu)} \varepsilon_{x} \\
\sigma_{y} & =\frac{\nu E}{(1+v)(1-2 \nu)} e+\frac{E}{(1+\nu)} \varepsilon_{y} \\
\sigma_{z} & =\frac{v E}{(1+v)(1-2 \nu)} e+\frac{E}{(1+v)} \varepsilon_{z} \\
e & =\varepsilon_{x}+\varepsilon_{y}+\varepsilon_{z}
\end{aligned}
$$

## Stress - strain relations 2-D case

$$
\left.\begin{array}{rlrl}
\sigma_{x} & =\frac{E}{1-v^{2}}\left(\varepsilon_{x}+v \varepsilon_{y}\right) \\
\sigma_{y} & =\frac{E}{1-v^{2}}\left(\varepsilon_{y}+v \varepsilon_{x}\right)
\end{array}\right) \quad \begin{aligned}
& =\tau / G \\
&
\end{aligned}
$$

## Second moment of area (inertia)



$$
\begin{gathered}
I_{x x}=\int_{A} y^{2} \mathrm{~d} A=\int_{-d / 2}^{d / 2} b y^{2} \mathrm{~d} y=b\left[\frac{y^{3}}{3}\right]_{-d / 2}^{d / 2} \\
I_{x x}=\frac{b d^{3}}{12}
\end{gathered}
$$

## Second moment of area (inertia)



$$
\begin{gathered}
I_{x x}=\int_{A} y^{2} \mathrm{~d} A=\int_{-d / 2}^{d / 2} 2\left(\frac{d}{2} \cos \theta\right) y^{2} \mathrm{~d} y \\
I_{x x}=\int_{-\pi / 2}^{\pi / 2} d \cos \theta\left(\frac{d}{2} \sin \theta\right)^{2} \frac{d}{2} \cos \theta \mathrm{~d} \theta \\
I_{x x}=\frac{d^{4}}{8} \int_{-\pi / 2}^{\pi / 2} \cos ^{2} \theta \sin ^{2} \theta \mathrm{~d} \theta \\
I_{x x}=\frac{\pi d^{4}}{64}
\end{gathered}
$$

## Product second moment of area (inertia)




$$
I_{x y}=\int_{A} x y \mathrm{~d} A
$$

## Parallel axes theorem (Steiner principle)



$$
I_{\mathrm{N}}=I_{\mathrm{C}}+A b^{2}
$$

## Aproximations for Thin-Walled sections




$$
I_{x x}=2\left[\frac{(b+t / 2) t^{3}}{12}+\left(b+\frac{t}{2}\right) t h^{2}\right]+t \frac{[2(h-t / 2)]^{3}}{12}
$$

$$
\begin{gathered}
I_{x x}=2\left[\frac{(b+t / 2) t^{3}}{12}+\left(b+\frac{t}{2}\right) t h^{2}\right]+\frac{t}{12}\left[(2)^{3}\left(h^{3}-3 h^{2} \frac{t}{2}+3 h \frac{t^{2}}{4}-\frac{t^{3}}{8}\right)\right] \\
I_{x x}=2 b t h^{2}+t \frac{(2 h)^{3}}{12}
\end{gathered}
$$

## Aproximations for Thin-Walled sections



$$
\begin{gathered}
I_{x x}=\int_{0}^{\pi r} t y^{2} \mathrm{~d} s \\
I_{x x}=\int_{0}^{\pi} t(r \cos \theta)^{2} r \mathrm{~d} \theta \\
I_{x x}=\frac{\pi r^{3} t}{2}
\end{gathered}
$$

## Thin-Walled sections - inclined walls



$$
I_{x x}=2 \int_{0}^{a / 2} t y^{2} \mathrm{~d} s=2 \int_{0}^{a / 2} t(s \sin \beta)^{2} \mathrm{~d} s
$$

$$
\begin{aligned}
& I_{x x}=\frac{a^{3} t \sin ^{2} \beta}{12} \\
& I_{y y}=\frac{a^{3} t \cos ^{2} \beta}{12} \\
& I_{x y}=\frac{a^{3} t \sin 2 \beta}{24}
\end{aligned}
$$

